

# Magnetic proximity interactions in the CrCl<sub>3</sub>/hBN/WSe<sub>2</sub> heterostructure

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Monolayers (MLs) of WSe<sub>2</sub> are part of the semiconducting transition metal dichalcogenide (S-TMD) family, which low-temperature ( $T=5-10$  K) photoluminescence (PL) spectra are composed of a series of excitonic lines, which have been widely investigated in the literature [1]. On the other hand, chromium chloride belonging to the family of magnetic layered materials is characterised by the ferromagnetic order with layers and the antiferromagnetic coupling between the layers with the in-plane spin alignment [2]. The combination of these two materials allows the study of the proximity effect [3].

In this work, we performed the photoluminescence (PL) experiment on the CrCl<sub>3</sub>/hBN/WSe<sub>2</sub> heterostructure (HS), composed of a thick CrCl<sub>3</sub> layer, hBN spacer and WSe<sub>2</sub> ML. The spacer plays a crucial role in preventing charge transfer from WSe<sub>2</sub> to CrCl<sub>3</sub>, as we presented in our previous work [3]. Measurements were performed at low temperature ( $T = 5$  K), using 2.41 eV excitation energy in an out-of-plane magnetic field.

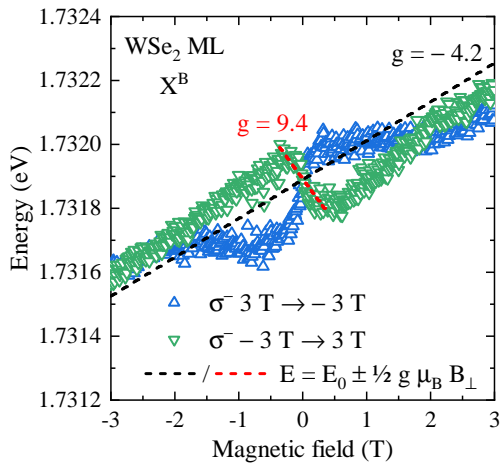


Figure 1: The magnetic-field evolution of the neutral exciton ( $X^B$ ) measured in the HS. The black line represents the Zeeman splitting function for the WSe<sub>2</sub> ML embedded in hBN flake based on the literature [1], while red dashed line is the fit using the Zeeman splitting equation.

We propose that the departure from the linear dependence of the  $X^B$  energy in external magnetic fields can be understood in terms of the interaction between the magnetic CrCl<sub>3</sub> and the WSe<sub>2</sub> ML. It is conceivable that the unusual energy dependence at small magnetic fields is due to the presence of the CrCl<sub>2</sub> magnetisation, while the increase in external magnetic fields causes the disappearance of the coupling between CrCl<sub>3</sub> and WSe<sub>2</sub>. Finally, the change in the orientation of the spin alignments in CrCl<sub>3</sub> above the critical fields results in the standard Zeeman splitting of the  $X^B$  line. Our results emphasise the complexity of the magnetic coupling between magnetically ordered CrCl<sub>3</sub> and the WSe<sub>2</sub> ML due to the proximity effect, which understanding is key for future possible applications.

The Figure represents the energy dependence of the neutral exciton ( $X^B$ ) measured on the HS in magnetic fields. Green and blue triangles denote the  $\sigma^-$  polarisation of the magnetic field course from 3 T to -3 T and from -3 T to 3 T, respectively. Within the range of small magnetic fields, an unusual behaviour of the  $X^B$  transition is observed. The energy values exhibit non-linear dependence as a function of magnetic fields. Around  $|0.2|$  T, a sudden change in energy of approximately  $3 \mu\text{eV}$  is observed. Using the standard equation for the Zeeman effect, describing the splitting into two circularly polarised components of the transition in the magnetic field [1], *i.e.*  $E^{\sigma^\pm}(B) = E_0 \pm 1/2 g \mu_B B_\perp$ , we found a  $g$ -factor of around 9.4 in this region. Furthermore, in magnetic field ranges of  $|0.2|$  T -  $|2|$  T and  $|2|$  T -  $|3|$  T, the data points form linear lines, each with different slopes. The magnetic field dependence of the  $X^B$  emission in the WSe<sub>2</sub> ML encapsulated in hBN flakes is characterised by the  $g$ -factor of around -4.2, which is represented by the dashed black line in the Figure. At the same time, the PL intensity of CrCl<sub>3</sub> shows a minimum at around  $|0.1|$  T. This is followed by the flatter of its intensity above  $|2|$  T, ascribed to a critical field [4], due to the change of the spin orientation in CrCl<sub>3</sub> from the intrinsic in-plane to along the applied out-of-plane magnetic field.

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[2] M. Gilbertini, *et al.* *Nat. Nanotechnol.* **14**, 408-419 (2019).

[3] Ł. Kipczak, *et al.* *arXiv:2304.11896* (2023).

[4] X. Cai, *et al.* *Nano Lett.* **19**, 3993-3998 (2019).